

Exercise / We derive the first-order energy correction in nondegenerate perturbation theory.

(Source: Griffiths 7.1).

Suppose we have some unperturbed system which we can solve exactly:

$$H^0 |\Psi_n^0\rangle = E_n^0 |\Psi_n^0\rangle, \quad \langle \Psi_m^0 | \Psi_n^0 \rangle = \delta_{nm}.$$

Next, we introduce a perturbed Hamiltonian which we cannot solve exactly:

$$H = H^0 + \lambda H'.$$

We Taylor expand the states and energies

$$|\Psi_n\rangle \approx |\Psi_n^0\rangle + \lambda |\Psi_n^1\rangle$$

$$E_n \approx E_n^0 + \lambda E_n^1,$$

and then we can expand the perturbed Schrödinger equation

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle \Rightarrow (H^0 + \lambda H') (|\Psi_n^0\rangle + \lambda |\Psi_n^1\rangle)$$

$$= (E_n^0 + \lambda E_n^1) (|\Psi_n^0\rangle + \lambda |\Psi_n^1\rangle) \Rightarrow H^0 |\Psi_n^0\rangle + \lambda H^0 |\Psi_n^1\rangle$$

$$+ \lambda H' |\Psi_n^0\rangle = E_n^0 |\Psi_n^0\rangle + \lambda E_n^0 |\Psi_n^1\rangle + \lambda E_n^1 |\Psi_n^0\rangle,$$

where we have only kept terms linear in λ . Grouping terms by λ , we have two equations:

$$H^0 |\Psi_n^0\rangle = E_n^0 |\Psi_n^0\rangle, \text{ our unperturbed system, and}$$

$$H^0 |\Psi_n^1\rangle + H' |\Psi_n^0\rangle = E_n^0 |\Psi_n^1\rangle + E_n^1 |\Psi_n^0\rangle \implies$$

$$\langle \Psi_n^0 | H^0 | \Psi_n^1 \rangle + \langle \Psi_n^0 | H' | \Psi_n^0 \rangle = E_n^0 \langle \Psi_n^0 | \Psi_n^1 \rangle +$$

$$E_n^1 \langle \Psi_n^0 | \Psi_n^0 \rangle \implies E_n^0 \langle \Psi_n^0 | \Psi_n^1 \rangle + \langle \Psi_n^0 | H' | \Psi_n^0 \rangle$$

$$= E_n^0 \langle \Psi_n^0 | \Psi_n^1 \rangle + E_n^1 \implies$$

$$\boxed{E_n^1 = \langle \Psi_n^0 | H' | \Psi_n^0 \rangle}$$

That is, the first correction to the n^{th} energy level is given by the expectation of the perturbation in the n^{th} unperturbed state.